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Driven vortices in confined geometry: the Corbino disk

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The fabrication of artificial pinning structures allows a new generation of experiments which can probe the properties of vortex arrays by forcing them to flow in confined geometries. We discuss the theoretical analysis of such experiments in both flux liquids and flux solids, focusing on the Corbino disk geometry. In the liquid, these experiments can probe the critical behavior near a continuous liquid-glass transition. In the solid, they probe directly the onset of plasticity.

1. INTRODUCTION

In the mixed state of type-II superconductors the magnetic field is concentrated in an array of flexible flux bundles that, much like ordinary matter, can form crystalline, liquid and glassy phases.[1,2] In clean systems the vortex solid melts into a flux liquid via a first order phase transition.[1] If the barriers to vortex line crossing are high, a rapidly cooled vortex liquid can bypass the crystal phase and get trapped in a metastable polymer-like glass phase.[3] The diversity of vortex structures is further increased by pinning from material disorder, which leads to a variety of novel glasses. Disorder-driven glass transitions are continuous, with diverging correlation lengths and universal critical behavior.[4,5]

Of particular interest is the dynamics of the vortex array in the various phases and in the proximity of a phase transition. In the liquid phase the vortex array flows yielding a linear resistivity. In clean systems the vortex solid melts into a flux liquid via a first order phase transition.[1] If the barriers to vortex line crossing are high, a rapidly cooled vortex liquid can bypass the crystal phase and get trapped in a metastable polymer-like glass phase.[3] The correlation length controlling the nonlocality of the flow grows with the liquid shear viscosity, which becomes large as the liquid freezes. At a continuous liquid-glass transition this correlation length diverges with a universal critical exponent. In the solid phase the vortex array moves as a single elastic object under uniform drive, provided the shear stresses are not too large. In the presence of strong spatial inhomogeneities, plastic flow occurs for large drives (or even for vanishingly small drives in a glassy solid) and the response is always nonlinear.[8] The dynamical correlation length can be identified with the separation between free dislocations and diverges at a continuous melting transition. Probing spatial velocity correlations can therefore give information on vortex dynamics within a given phase, as well as on the nature of the phase transitions connecting the various phases.

As for ordinary matter, the shear rigidity of the vortex array can be probed by forcing the vortices to flow in confined geometries.[6,9] This type of experiments was pioneered by Kes and collaborators to study the shear rigidity of the two-dimensional vortex liquid near freezing in thin films.[10] More recently, patterned irradiation of cuprate superconductors with heavy ions has made it possible to create samples with controlled distributions of damage tracks.[11] We recently showed that an analysis of such experiments that combines an inhomogeneous scaling theory with the hydrodynamics of viscous flux liquids can be used to infer the critical behavior near a continuous glass transition, as well as to distinguish between continuous transitions, such as that to a Bose glass, and nonequilibrium transition to a polymer-like glass driven by interaction.

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2. LIQUID FLOW IN CHANNELS

In the Corbino disk, with magnetic field along the disk axis (z direction), a uniform radial current density of magnitude \( J(r) = I/(2\pi r) \) is introduced in the sample by injecting current at the center and removing it at the outer circumference of the disk (inset of Fig. 1). The current drives the vortices to move in circles about the axis. In the flux liquid, the dynamics on scales larger than the intervortex spacing, \( a_0 \), is described by hydrodynamic equations for the flow velocity \( \mathbf{v}(r) \), which determines the local field from flux motion,

\[
\mathbf{E} = n_0 \phi_0 \hat{\mathbf{z}} \times \mathbf{v}(r)/c, \quad \text{with} \quad n_0 = 1/a_0^2.
\]

For simple geometries like the Corbino disk, where the current is spatially homogeneous in the \( z \) direction, hydrodynamics reduces to a single equation,

\[
- \gamma \mathbf{v} + \eta \nabla^2 \mathbf{v} = \frac{1}{c} n_0 \phi_0 \hat{\mathbf{z}} \times \mathbf{J}(r),
\]

where \( \gamma(T, H) \) is the friction, \( \eta(T, H) \) is the viscosity controlling the viscous drag from interactions and entanglement, and the term on the right hand side is the Lorentz force density driving flux motion. It is instructive to rewrite Eq. (1) as an equation for the local field,

\[
- \xi^2 \nabla^2 \mathbf{E} + \mathbf{E} = \rho_f \mathbf{J},
\]

with \( \xi = \sqrt{\eta/\gamma} \) the viscous correlation length and \( \rho_f = (n_0 \phi_0/c)^2 \gamma \) the flux flow resistivity. If the viscous force is negligible, Eq. (2) is simply Ohm’s law and the radial field is \( E_0(r) = (\rho_f/2\pi t)(1/r) \).

To probe the viscous drag, it is necessary to force large scale spatial inhomogeneities in the flow. This may be achieved by suitable pinning boundaries. As an example, we imagine selectively irradiating a cylindrical central region and an outer annular region of the disk to obtain the structure sketched in the inset of Fig. 2. Here the vortices in the heavily irradiated central and outer regions (shaded) are in the Bose glass phase, while vortices in the unirradiated (white) annular region are in the flux liquid phase. A radial current drives tangential flow in the resistive flux liquid annulus, which is impeded by the “Bose-glass contacts” at the boundaries. The field profile obtained by solving Eq. (2) with no-slip boundary conditions is spatially inhomogeneous on length \( \xi \), as shown in Fig. 1. One can probe this profile and extract \( \xi \) by placing a string of radial contacts at \( r_n \), for \( n = 1, 2, 3, \ldots \), and measuring the voltage \( V_{n+1,n} \) across each successive pair (inset of Fig. 2). If the viscosity is small \( (\xi << d) \), the voltage decreases logarithmically as one moves from the inner to the outer contacts, as in a freely flowing uncorrelated liquid, where \( V_{n+1,n}^0 = (\rho_f/2\pi t) \ln(r_{n+1}/r_n) \). When \( \xi \) grows, the onset of rigidity in the liquid becomes apparent (Fig. 2). An elastic vortex solid would ro-
The voltage drop $2\pi V_{n+1,n}/(\rho_f l)$ across pairs of contacts $(r_{n+1}, r_n)$, with $r_n = R_1 + nd$, for $n = 0, 2, \ldots, 10$, $R_1 = d$, and $d = W/10$ the contact spacing. The symbols refer to $\xi/d = 0.1$ (triangles), $\xi/d = 1$ (squares) and $\xi/d = 2$ (circles). Solid lines are guides to the eye. Inset: top view of the Corbino disk with Bose glass contacts, with $W = R_2 - R_1$.

![Diagram of voltage drop across pairs of contacts](image)

**Table 1**

<table>
<thead>
<tr>
<th>Transition</th>
<th>$\gamma(T)$</th>
<th>$\eta(T)$</th>
</tr>
</thead>
<tbody>
<tr>
<td>Bose Glass</td>
<td>$t^{(1-z)\nu_\perp}$</td>
<td>$t^{-(z+1)\nu_\perp}$</td>
</tr>
<tr>
<td>Vortex Glass</td>
<td>$t^{(2-z)\nu_\perp}$</td>
<td>$t^{-z\nu_\perp}$</td>
</tr>
<tr>
<td>“Polymer” Glass</td>
<td>finite</td>
<td>$e^{n/t}$</td>
</tr>
<tr>
<td>First order freezing</td>
<td>finite</td>
<td>jumps to $\infty$</td>
</tr>
<tr>
<td>Continuous freezing</td>
<td>finite</td>
<td>$e^{c/t^{0.369\ldots}}$</td>
</tr>
</tbody>
</table>

Behavior of $\gamma$ and $\eta$ near various transitions of vortex matter. Here $t = (T - T_0)/T_0$, with $T_0$ the relevant transition temperature in each case.

...cosity, $\eta$. Experiments in artificial pinning structures can therefore be used to infer both $\gamma$ and $\eta$. The behavior of these two parameters near a phase transition can in fact be used to identify the transition, as summarized in table 1.

Of particular interest is the case of continuous glass transitions, which are characterized by universal critical behavior. Experiments of the type just described are especially powerful in this case as they can be used to map out the critical behavior. One can probe the Bose glass transition by lightly irradiating the liquid annular region, and then lowering the temperature at constant field from $T_{\text{annulus}}^\text{BG} < T < T_{\text{contacts}}^{\text{BG}}$. Most physical properties near the transition can be described via a scaling theory in terms of diverging correlation lengths perpendicular and parallel to the field direction, $\xi_\perp(T) \sim |T - T_{\text{BG}}^\text{annulus}|^{-\nu_\perp}$ and $\xi_\parallel(T) \sim |T - T_{\text{BG}}^\text{annulus}|^{-\nu_\parallel}$, with $\nu_\parallel = 2\nu_\perp$, and a diverging correlation time, $\tau \sim \xi_\perp^{z_\perp} \sim |T - T_{\text{BG}}^\text{annulus}|^{-z\nu_\perp}$. Scaling can then be used to relate physical quantities to these diverging length and time scales. In particular, the friction coefficient $\gamma$ that determines the bulk flux flow resistivity $\rho_f(T)$ is predicted to diverge as $T \to T_{\text{BG}}^\text{annulus}$ as $\rho_f \sim |T - T_{\text{BG}}^\text{annulus}|^{-\nu_\perp}(z - 2)$. As shown recently by Marchetti and Nelson, when flux flow in confined geometries is analyzed by combining hydrodynamics with the the Bose glass scaling theory – generalized to the spatially inhomogeneous case – the Bose glass correlation length $\xi_\perp$ is naturally identified with the viscous length $\xi_\parallel$. It then follows that the liquid shear viscosity diverges at $T_{\text{BG}}$ as $\eta \sim |T - T_{\text{BG}}|^{-z\nu_\perp}$. Furthermore, the scaling of the finite-geometry...
resistivity displayed in Eq. 3 is a general property of the vortex liquid near a continuous glass transition. Hydrodynamics yields the precise form of the scaling function, which depends on the experimental geometry and can be found in Ref. 3 for the Corbino disk.

3. PLASTIC FLOW IN DRIVEN SOLIDS

As shown in recent experiments by the Argonne group, the Corbino disk geometry can also be used to study the onset of plastic flow in a driven solid. In this case we consider an unirradiated disk, where the vortex array has a clear melting transition. Below melting the vortex solid moves as a rigid body, with \( v(r) \sim r \), and the voltage grows as \( r^2 \). The \( \sim 1/r \) dependence of the driving force yields, however, large elastic deformations of the medium, described by the solution of

\[
\sigma_{\phi\phi}(r) = \left( n_0 \phi_0 I / 4 \pi c t \right) \left[ 1 + 2 \ln \left( R_2 / R_1 \right) R_1^2 / r^2 \right],
\]

which the entire vortex solid shear melts. The melting radius \( R_M \) increases with current, indicating that, since the stresses are largest near the axis of the disk, “shear-induced melting” occurs first in circular layers close to the axis. This behavior is qualitatively consistent with the observations by the Argonne group. The simple model described here suggests that at high fields the current scale \( I_0 \) is independent of field. A more detailed calculation incorporating field and temperature dependence will be described elsewhere.

REFERENCES