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# Gluinos condensing at the CCNI: 4096 CPUs weigh in\*

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## Abstract

We report preliminary results of lattice super-Yang-Mills computations using domain wall fermions, performed at an actual rate of 1000 Gflop/s, over the course of six months, using two BlueGene/L racks at Rensselaer's CCNI supercomputing center. This has allowed us to compute the gluino condensate and string tension over a wide range of lattice parameters, setting the stage for continuum, chiral extrapolations.

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In this talk, I present preliminary results obtained in collaboration with coauthors of a forthcoming paper [1]. We have used domain wall fermions (DWF) [2, 3] to study nonperturbative aspects of pure  $\mathcal{N} = 1$  super-Yang-Mills (SYM) [4]. The strong dynamics of supersymmetric gauge theories underlie most models of spontaneous supersymmetry (SUSY) breaking, and the development of a first-principles tool is needed. Though at finite lattice spacing SUSY is violated, it is automatically recovered [5] in the continuum limit with a massless gluino.<sup>1</sup> When the DWF formulation is employed, this “chiral limit” is achieved without the need for a computationally expensive, nonperturbatively determined fine-tuning of the bare gluino mass [6, 7]; in the limit of infinite domain wall separation  $L_s \rightarrow \infty$ , DWF realize the lattice chiral symmetry [8] associated with Ginsparg-Wilson fermions [9], which protects against additive mass renormalization. These nice features of the DWF approach are to be contrasted with the Wilson fermion formulation, which was pursued for several years by the DESY-Münster-Roma collaboration [10, 11, 12, 13, 14, 15, 16].

Lattice studies can provide details that other approaches cannot, such as “snapshots” of the gauge field configurations that are dominating the gluino condensate. For instance, in the work of Fleming et al. [17], the only DWF simulation of SYM to date, it was suggested that spikes in the gluino condensate may correspond to configurations with fractional topological charge, as would be expected from Ref. [18]. Further pursuit of this conjecture, consisting of lattice studies of monopoles and topological charge, is on our agenda. Also, we will compute the low-lying spectrum of composite states, consisting of strongly bound gluons and gluinos. Apart from the inconclusive results of DESY-Münster-Roma, this aspect of SYM is completely unknown from continuum methods, and ideally suited to the lattice approach.

At early stages in such studies, understanding numerical behavior of important quantities such as the gluino condensate will teach us a lot about the lattice formulation that we currently do not know. For instance, it is important to set “benchmarks” regarding compute time, and lattice artifacts such as discretization and finite size effects. The first place that we will examine this is in the gluino condensate, which is believed to be known exactly by continuum methods [18, 19, 20, 21, 22], and is therefore ideal for calibrating the lattice methods. Understanding of the lattice theory and simulation performance is already emerging from our preliminary results, as we now briefly discuss.

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<sup>1</sup> Here we use the terms “gluon” and “gluino” by way of analogy. It should be kept in mind that the strongly coupled gauge theory would be an extension to the gauge group of the Standard Model.

Domain wall fermion simulations require world-class computing resources, such as are available to Giedt at Rensselaer; namely, the Computational Center for Nanotechnology Innovations (CCNI), one of the world’s most powerful university-based supercomputing centers, and a top 25 supercomputing center of any kind in the world. We are presently the third heaviest user of this facility, and have been generating lattice configurations and measurements continuously at a sustained actual rate of 1000 Gflop/s since the end of January 2008. For comparison the DESY-Münster-Roma collaboration performed their computations at a rate of 10 Gflop/s for a cumulative time of one year. Thus our study represents a hundred-fold improvement over what has been done previously, just in terms of raw computation power.

As an example of our results, we have obtained the bare gluino condensate from dynamical domain wall fermion simulations for a variety of bare gauge couplings  $g$ , parameterized in terms of  $\beta = 4/g^2$ , as is conventional in SU(2) lattice gauge theory. The results for a  $16^3 \times 32$  lattice (i.e., the number of sites in spatial and temporal directions) with domain wall separation  $L_s = 16$  sites are displayed in Fig. 1. We note that such data for the condensate versus  $\beta$  has never been obtained before; it is important because the continuum limit corresponds to  $\beta \rightarrow \infty$  (with the physical size of the lattice held fixed). A fit to the data obviously yields a vanishing condensate at a finite gauge coupling  $\beta \sim 2.7$ . This just reflects the fact that as  $\beta$  increases the lattice spacing shrinks, and thus so does the physical size of the lattice in its entirety. In a small enough “box” confinement will disappear and the condensate “melts.” Thus we already gain an important benchmark: to go much beyond  $\beta = 2.5$  will require larger lattices, and in fact one should carefully measure systematic errors due to finite size effects at  $\beta \approx 2.5$ . This is consistent with what is already known from the so-called “quenched” theory, which has no gluinos.

Fig. 2 shows the gluino condensate for decreasing values of the residual mass  $m_{\text{res}}$ , which is a measure of explicit chiral symmetry breaking due to finite  $L_s$  [23]. As expected, larger  $L_s$  values have the smallest  $m_{\text{res}}$ , and a nonzero gluino condensate appears to occur in the  $m_{\text{res}} \rightarrow \infty$  limit. Also as expected, smaller values of  $m_{\text{res}}$  occur for the weaker coupling  $\beta = 2.4$ .

Finally, we have looked at Creutz ratios [24],

$$\chi(R, R) = -\ln \frac{W(R, R)W(R-1, R-1)}{W(R, R-1)^2} \sim \sigma a^2, \quad (0.1)$$

where  $W(R, R')$  is an  $R \times R'$  Wilson loop, in order to extract the string tension  $\sigma$  in lattice units, as well as to delineate the scaling regime where the continuum limit may be extracted. In the process we obtain an estimate of the lattice spacing in

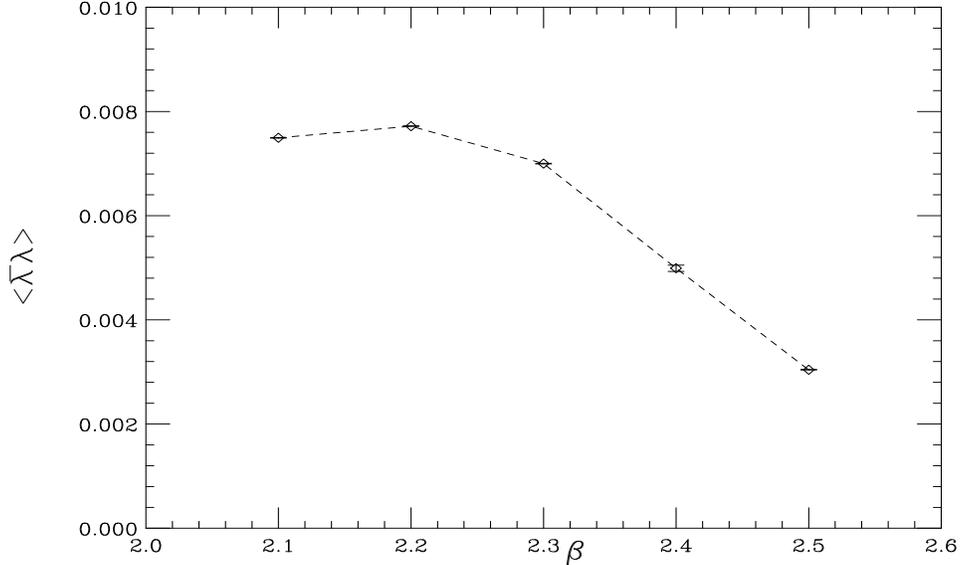


FIG. 1: The gluino condensate versus  $\beta$  for a  $16^3 \times 32$  lattice with domain wall separation  $L_s = 16$  (dashed line drawn to guide the eye). It can be estimated from the figure that for the  $16^3 \times 32$  lattice studied here, the system will deconfine at  $\beta \sim 2.7$ , as a result of finite size effects.

units of the string tension. Results for the  $16^3 \times 32$ ,  $L_s = 16$  lattice are shown in Fig. 3. Although the errors are somewhat large, scaling is clearly setting in at around  $\beta = 2.3$ . To see this one notes that the larger Creutz ratios appear to coalesce on an envelope, corresponding to the distance scale at which an area law begins to take hold in the Wilson loops.

We emphasize that all of the results mentioned above are ground-breaking as far as lattice SYM is concerned. Several important goals related to the lattice SYM project have already been achieved, laying the groundwork for the more extensive studies that will follow:

- Developed parallel simulation code for SYM by modification of the current version of the Columbia Physics System (CPS) QCD package.
- This extends DOE funded code (CPS, part of USQCD’s SciDAC program) to “beyond the Standard Model” physics, which is a realization of one of the USQCD Collaboration objectives [25].
- We have reproduced the results of [17] as a check on our code.

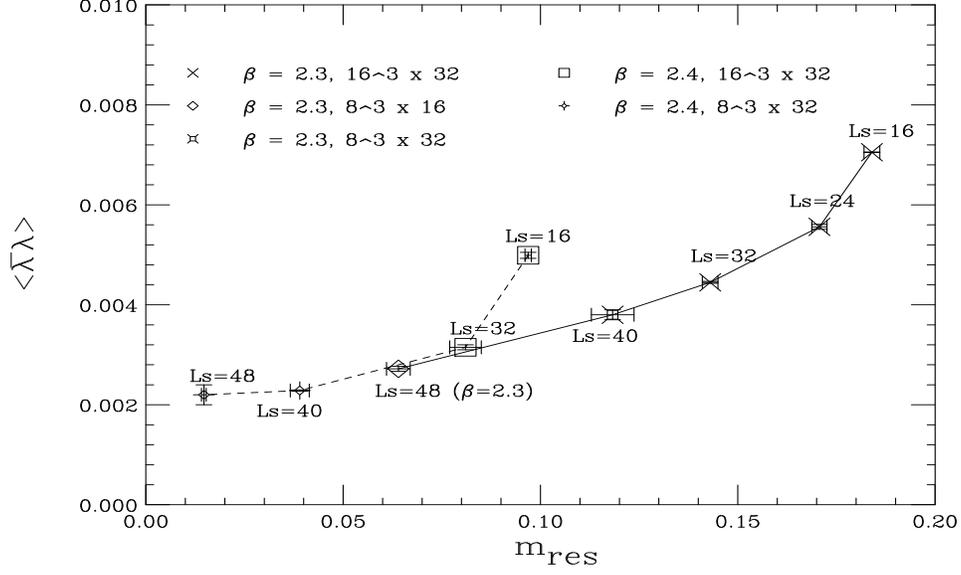


FIG. 2: Our simulation results for the gluino condensate versus  $m_{\text{res}}$ , where the latter is a measure of explicit chiral symmetry breaking due to finite domain wall separation  $L_s$ . The solid line corresponds to  $\beta = 2.3$  whereas the dashed line is for  $\beta = 2.4$ .

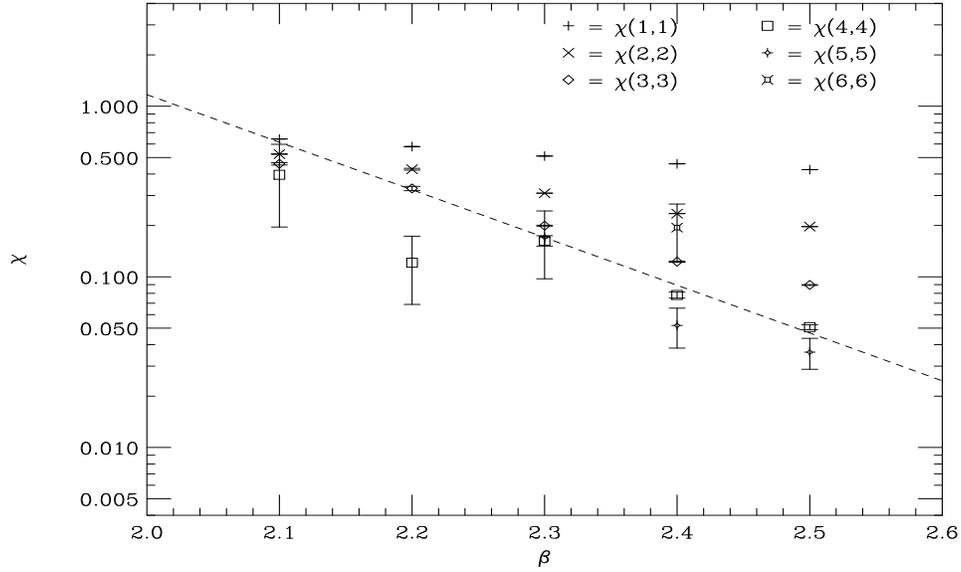


FIG. 3: Creutz ratios for the  $16^3 \times 32$  lattice with  $L_s = 16$ . The dashed line indicates the 2-loop SUSY prediction for the dependence  $\chi \sim \sigma a^2$  where  $\sigma$  is the string tension and  $a = a(\beta)$  is the lattice spacing.

- Our software runs successfully on IBM’s Bluegene (BG) architecture, taking full advantage of BG specific communications utilities.
- Developed Landau gauge-fixing and Fourier space propagator code for adjoint fermion representations, essential for nonperturbative renormalization of the condensate, in the RI/MOM scheme [26, 27].
- Established timing and statistical uncertainty benchmarks. For example, we have found that for small single-node volumes ( $16^3 \times 32 \times L_s/2048$  CPU’s =  $64 \times L_s$  sites per CPU) the efficiency of the BG parallel code is 10 percent.

In the course of our studies, we intend to investigate other domain wall fermion formulations (e.g., “gap” [28] and “Mobius” [29]) as ways to approach the chiral limit more quickly. Also, we will implement recent optimizations of fermion matrix inverters (to increase efficiency) and improved actions (to reduce systematic errors).

### Acknowledgements

We benefited from a copy of the code that was used in [17], and employed it as the basis for our modifications to the current version of the Columbia Physics System. At various points JG benefited from technical assistance provided by Chulwoo Jung (Brookhaven National Lab) and Adam Todorski (SCOREC and CCNI at Rensselaer). The computational efforts on this project mainly utilized the Computational Center for Nanotechnology Innovations (CCNI), and JG expresses his appreciation for continuous access to that facility. The project also, at times, utilized the SUR BlueGene/L at Rensselaer, which is supported by NSF grant 0420703 entitled “MRI: Acquisition of Infrastructure for Research in Grid Computing and Multiscale Systems Computation” and a gift by the IBM Corporation of a BlueGene/L computer. JG acknowledges support from Rensselaer faculty development funds.

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